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Supplementary Release on Backward Equations for the Functions $T(p,h)$, $v(p,h)$ and $T(p,s)$, $v(p,s)$ for Region 3 of the IAPWS Industrial Formulation 1997 for the Thermodynamic Properties of Water and Steam

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President:

Professor Koichi Watanabe
Department of System Design Engineering
Keio University
3-14-1 Hiyoshi, Kohoku-ku
Yokohama 223-8522, Japan

Executive Secretary:

Dr. R. B. Dooley
Electric Power Research Institute
3412 Hillview Avenue
Palo Alto, California 94304, USA

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The backward equations for temperature and specific volume as functions of pressure and enthalpy $T(p,h)$, $v(p,h)$ and as functions of pressure and entropy $T(p,s)$, $v(p,s)$ for region 3 provided in this release are recommended as a supplement to "IAPWS Industrial Formulation 1997 for the Thermodynamic Properties of Water and Steam" (IAPWS-IF97) [1, 2] and to "Backward Equations for Pressure as a Function of Enthalpy and Entropy $p(h,s)$ to the Industrial Formulation 1997 for the Thermodynamic Properties of Water and Steam" (referred to here as IAPWS-IF97-S01) [3, 4]. Further details concerning the equations can be found in the corresponding article by H.-J. Kretzschmar et al. [5].

Further information concerning this supplementary release, IAPWS-IF97, IAPWS-IF97-S01, and other releases issued by IAPWS can be obtained from the Executive Secretary of IAPWS or from <http://www.iapws.org>.

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1 Nomenclature

Thermodynamic quantities:

| | |
|----------|-----------------------------------|
| f | Specific Helmholtz free energy |
| h | Specific enthalpy |
| p | Pressure |
| s | Specific entropy |
| T | Absolute temperature ^a |
| v | Specific volume |
| Δ | Difference in any quantity |
| h | Reduced enthalpy, $h = h/h^*$ |
| q | Reduced temperature $q = T/T^*$ |
| p | Reduced pressure, $p = p/p^*$ |
| r | Density |
| s | Reduced entropy, $s = s/s^*$ |
| w | Reduced volume, $w = v/v^*$ |

Root-mean-square value:

$$\Delta x_{\text{RMS}} = \sqrt{\frac{1}{N} \sum_{n=1}^N (\Delta x_n)^2}$$

where Δx_n can be either absolute or percentage difference between the corresponding quantities x ; N is the number of Δx_n values (100 million points uniformly distributed over the range of validity in the p - T plane).

Superscripts:

| | |
|----|------------------------------------|
| 01 | Equation of IAPWS-IF97-S01 |
| 97 | Quantity or equation of IAPWS-IF97 |
| * | Reducing quantity |
| ' | Saturated liquid state |
| " | Saturated vapor state |

Subscripts:

| | |
|-----|---------------------------------------|
| 1 | Region 1 |
| 2 | Region 2 |
| 3 | Region 3 |
| 3a | Subregion 3a |
| 3b | Subregion 3b |
| 3ab | Boundary between subregions 3a and 3b |
| 4 | Region 4 |
| 5 | Region 5 |
| B23 | Boundary between regions 2 and 3 |
| c | Critical point |
| it | Iterated quantity |
| max | Maximum value of a quantity |
| RMS | Root-mean-square value of a quantity |
| sat | Saturation state |
| tol | Tolerated value of a quantity |

^a Note: T denotes absolute temperature on the International Temperature Scale of 1990 (ITS-90).

2 Background

The Industrial Formulation IAPWS-IF97 for the thermodynamic properties of water and steam [1, 2] contains basic equations, saturation equations and equations for the most often used backward functions $T(p, h)$ and $T(p, s)$ valid in the liquid region 1 and the vapor region 2; see Figure 1. IAPWS-IF97 was supplemented by "Backward Equations for Pressure as a Function of Enthalpy and Entropy $p(h, s)$ to the Industrial Formulation 1997 for the Thermodynamic Properties of Water and Steam" [3, 4], which we will refer to as IAPWS-IF97-S01, including equations for the backward function $p(h, s)$ valid in region 1 and region 2.

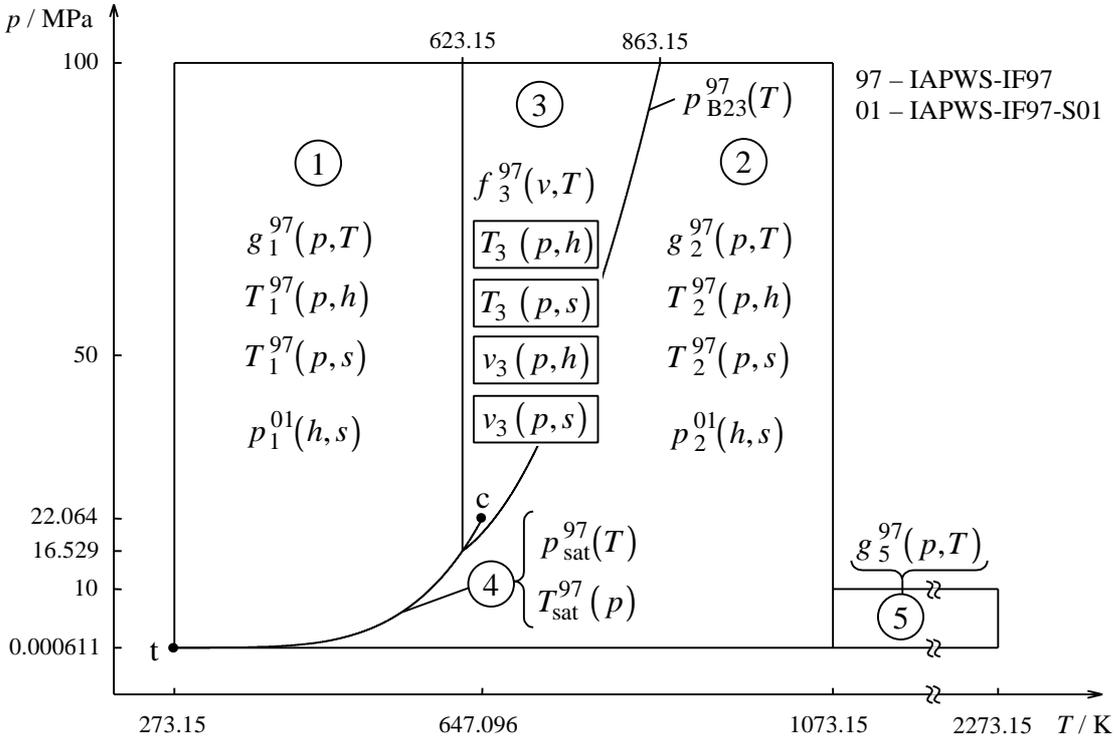


Figure 1. Regions and equations of IAPWS-IF97, IAPWS-IF97-S01, and the backward equations $T(p, h)$, $v(p, h)$, and $T(p, s)$, $v(p, s)$ of this release

In modeling steam power cycles, thermodynamic properties as functions of the variables (p, h) or (p, s) are also required in region 3. It is difficult to perform these calculations with IAPWS-IF97, because they require two-dimensional iterations using the functions $p(v, T)$, $h(v, T)$ or $p(v, T)$, $s(v, T)$ that can be explicitly calculated from the fundamental region 3 equation $f(v, T)$. While these calculations are not frequently required in region 3, the relatively large computing time required for two-dimensional iteration can be significant in process modeling.

In order to avoid such iterations, this release provides equations for the backward functions $T_3(p, h)$, $v_3(p, h)$ and $T_3(p, s)$, $v_3(p, s)$, see Figure 1. With temperature and specific

volume calculated from the backward equations, the other properties in region 3 can be calculated using the IAPWS-IF97 basic equation $f_3^{97}(v, T)$.

The numerical consistency with the IAPWS-IF97 basic equation $f_3^{97}(v, T)$ of T and v calculated from these backward equations is sufficient for most applications in heat cycle and steam turbine calculations. For applications where the demands on numerical consistency are extremely high, iterations using the IAPWS-IF97 basic equation $f_3^{97}(v, T)$ may be necessary. In these cases, the equations $T_3(p, h)$, $v_3(p, h)$ and $T_3(p, s)$, $v_3(p, s)$ can be used for calculating very accurate starting values.

The backward equations $T_3(p, h)$, $v_3(p, h)$ and $T_3(p, s)$, $v_3(p, s)$ can only be used in their ranges of validity described in Section 4. They should not be used for determining any thermodynamic derivatives.

In any case, depending on the application, a conscious decision is required whether to use the backward equations $T_3(p, h)$, $v_3(p, h)$ and $T_3(p, s)$, $v_3(p, s)$ or to calculate the corresponding values by iterations from the basic equation of IAPWS-IF97.

3 Numerical Consistency Requirements

The permissible value for the numerical consistency $|\Delta T|_{\text{tol}} = 25 \text{ mK}$ of the backward functions $T_3(p, h)$ and $T_3(p, s)$ with the basic equation $f_3^{97}(v, T)$ was determined by IAPWS [6, 7] as a result of an international survey.

The permissible value Δv_{tol} for the numerical consistency for the equations $v_3(p, h)$ and $v_3(p, s)$ can be estimated from the total differentials

$$\Delta v_{\text{tol}} = \left(\frac{\partial v}{\partial T} \right)_h \Delta T_{\text{tol}} + \left(\frac{\partial v}{\partial h} \right)_T \Delta h_{\text{tol}} \quad \text{and} \quad \Delta v_{\text{tol}} = \left(\frac{\partial v}{\partial T} \right)_s \Delta T_{\text{tol}} + \left(\frac{\partial v}{\partial s} \right)_T \Delta s_{\text{tol}} ,$$

where $\left(\frac{\partial v}{\partial T} \right)_h$, $\left(\frac{\partial v}{\partial h} \right)_T$, $\left(\frac{\partial v}{\partial T} \right)_s$, and $\left(\frac{\partial v}{\partial s} \right)_T$ are derivatives [8] calculated from the IAPWS-IF97 basic equation and Δh_{tol} and Δs_{tol} are values determined by IAPWS for the adjacent region 1 and subregion 2c [9], see Table 1. The resulting permissible specific volume difference is $|\Delta v/v|_{\text{tol}} = 0.01 \%$ for both functions $v_3(p, h)$ and $v_3(p, s)$.

At the critical point $\left[T_c = 647.096 \text{ K}, v_c = 1/(322 \text{ kg m}^{-3}) \right]$, more stringent consistency requirements were arbitrarily set. These were $|\Delta T|_{\text{tol}} = 0.49 \text{ mK}$ and $|\Delta v/v|_{\text{tol}} = 0.0001 \%$.

Table 1. Numerical consistency values $|\Delta T|_{\text{tol}}$ of [6] required for $T_3(p, h)$ and $T_3(p, s)$, values $|\Delta h|_{\text{tol}}$, $|\Delta s|_{\text{tol}}$ of [9], and resulting tolerances $|\Delta v/v|_{\text{tol}}$ required for $v_3(p, h)$ and $v_3(p, s)$

| | $ \Delta T _{\text{tol}}$ | $ \Delta h _{\text{tol}}$ | $ \Delta s _{\text{tol}}$ | $ \Delta v/v _{\text{tol}}$ |
|----------------|---------------------------|---------------------------|--|-----------------------------|
| Region 3 | 25 mK | 80 J kg ⁻¹ | 0.1 J kg ⁻¹ K ⁻¹ | 0.01 % |
| Critical Point | 0.49 mK | - | - | 0.0001 % |

4 Structure of the Equation Set

The equation set consists of backward equations $T(p, h)$, $v(p, h)$ and $T(p, s)$, $v(p, s)$ for region 3.

Region 3 is defined by:

$$623.15 \text{ K} \leq T \leq 863.15 \text{ K} \text{ and } p_{\text{B23}}^{97}(T) \leq p \leq 100 \text{ MPa},$$

where p_{B23}^{97} represents the B23 equation of IAPWS-IF97. Figure 2 shows the way in which region 3 is divided into the two subregions 3a and 3b.

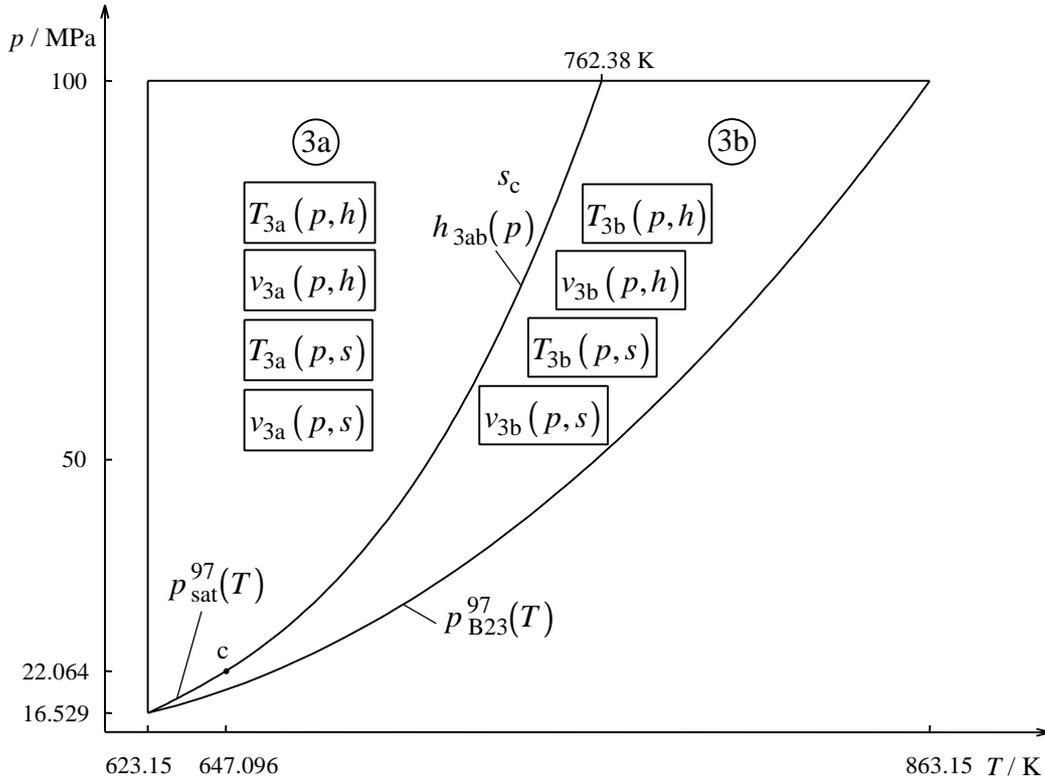


Figure 2. Division of region 3 into two subregions 3a and 3b for the backward equations $T(p, h)$, $v(p, h)$ and $T(p, s)$, $v(p, s)$

Table 2 shows the decisions which have to be made in order to find the correct subregion for the functions $T(p, h)$, $v(p, h)$ and $T(p, s)$, $v(p, s)$.

Table 2. Criteria for finding the correct subregion, 3a or 3b, for the backward functions $T(p, h)$, $v(p, h)$ and $T(p, s)$, $v(p, s)$

| Backward Functions $T(p, h)$, $v(p, h)$ | | | Backward Functions $T(p, s)$, $v(p, s)$ | | |
|--|---------------------|------------------|--|----------------|-----------------|
| | Subregion | | | Subregion | |
| | 3a | 3b | | 3a | 3b |
| for $p < p_c$: | $h \leq h'(p)$ | $h \geq h''(p)$ | for $p < p_c$: | $s \leq s'(p)$ | $s \geq s''(p)$ |
| for $p \geq p_c$: | $h \leq h_{3ab}(p)$ | $h > h_{3ab}(p)$ | for $p \geq p_c$: | $s \leq s_c$ | $s > s_c$ |

For pressures less than the critical pressure $p_c = 22.064$ MPa, the saturation line is the boundary between subregions 3a and 3b. That means for the functions $T(p, h)$ and $v(p, h)$, if the given specific enthalpy h is less than or equal to $h'(p)$ calculated from the given pressure p on the saturated liquid line, then the point of state to be calculated is located in subregion 3a. If the given enthalpy h is greater than or equal to $h''(p)$ calculated on the saturated vapor line, then the point of state is located in subregion 3b. Otherwise, the point is in the two-phase region. In that case, the saturation temperature equation $T_{\text{sat}}^{97}(p)$ and the basic equation $f_3^{97}(v, T)$ of IAPWS-IF97 can be used to calculate the temperature and the specific volume from the given pressure and the given enthalpy. The decisions are analogous for the functions $T(p, s)$ and $v(p, s)$.

For pressures greater than or equal to p_c , the boundary between the subregions 3a and 3b corresponds to the critical isentropic line $s = s_c$, see Figure 2. For the functions $T(p, s)$ and $v(p, s)$, input points can be tested directly to identify the subregion since the specific entropy is an independent variable. If the given specific entropy s is less than or equal to

$$s_c = 4.412\,021\,482\,234\,76 \text{ kJ kg}^{-1} \text{ K}^{-1},$$

then the state point to be calculated is located in subregion 3a; otherwise it is in subregion 3b. In order to decide which $T(p, h)$, $v(p, h)$ equation, 3a or 3b, must be used for given values of p and h , the boundary equation $h_{3ab}(p)$, Eq. (1), has to be used, see Figure 2. This equation is a polynomial of the third degree and reads

$$\frac{h_{3ab}(p)}{h^*} = \mathbf{h}(\mathbf{p}) = n_1 + n_2 \mathbf{p} + n_3 \mathbf{p}^2 + n_4 \mathbf{p}^3, \quad (1)$$

where $\mathbf{h} = h/h^*$ and $\mathbf{p} = p/p^*$ with $h^* = 1 \text{ kJ kg}^{-1}$ and $p^* = 1 \text{ MPa}$. The coefficients n_1 to n_4 of Eq. (1) are listed in Table 3. The range of the equation $h_{3ab}(p)$ is from the critical point to 100 MPa. The related temperature at 100 MPa is $T = 762.380\,873\,481 \text{ K}$. Equation (1) does not exactly describe the critical isentropic line. The maximum specific entropy deviation was determined as

$|\Delta s_{3ab}|_{\max} = \left| s_3^{97}(T_{\text{it}}^{97}(p, h_{3ab}(p)), v_{\text{it}}^{97}(p, h_{3ab}(p))) - s_c \right|_{\max} = 0.66 \text{ J kg}^{-1} \text{ K}^{-1}$,
 where T_{it}^{97} and v_{it}^{97} were obtained by iterations using the derivatives $p_3^{97}(v, T)$ and $s_3^{97}(v, T)$ of the IAPWS-IF97 basic equation for region 3.

Table 3. Numerical values of the coefficients of the equation $h_{3ab}(p)$ in its dimensionless form, Eq. (1), for defining the boundary between subregions 3a and 3b^a

| i | n_i | i | n_i |
|-----|---|-----|---|
| 1 | $0.201\ 464\ 004\ 206\ 875 \times 10^4$ | 3 | $-0.219\ 921\ 901\ 054\ 187 \times 10^{-1}$ |
| 2 | $0.374\ 696\ 550\ 136\ 983 \times 10^1$ | 4 | $0.875\ 131\ 686\ 009\ 950 \times 10^{-4}$ |

^a It is recommended that programmed functions be verified using 8 byte real values for all variables.

If the given specific enthalpy h is greater than $h_{3ab}(p)$ calculated from the given pressure p , then the state point to be calculated is located in subregion 3b, otherwise it is in subregion 3a (see Figure 2).

Note, Eq. (1) does not correctly simulate the isentropic line $s = s_c$ at pressures lower than p_c . However, the calculated values $h_{3ab}(p)$ are not higher than $h''(p)$ and not lower than $h'(p)$.

For *computer-program verification*, Eq. (1) gives the following p - h point:

$$p = 25 \text{ MPa}, \quad h_{3ab}(p) = 2.095\ 936\ 454 \times 10^3 \text{ kJ kg}^{-1}.$$

5 Backward Equations $T(p, h)$ and $v(p, h)$ for Subregions 3a and 3b

5.1 The Equations $T(p, h)$

The backward equation $T_{3a}(p, h)$ for subregion 3a has the following dimensionless form:

$$\frac{T_{3a}(p, h)}{T^*} = \mathbf{q}_{3a}(\mathbf{p}, \mathbf{h}) = \sum_{i=1}^{31} n_i (\mathbf{p} + 0.240)^{I_i} (\mathbf{h} - 0.615)^{J_i}, \quad (2)$$

where $\mathbf{q} = T/T^*$, $\mathbf{p} = p/p^*$, and $\mathbf{h} = h/h^*$ with $T^* = 760 \text{ K}$, $p^* = 100 \text{ MPa}$, and $h^* = 2300 \text{ kJ kg}^{-1}$. The coefficients n_i and exponents I_i and J_i of Eq. (2) are listed in Table 4.

The backward equation $T_{3b}(p, h)$ for subregion 3b reads in its dimensionless form

$$\frac{T_{3b}(p, h)}{T^*} = \mathbf{q}_{3b}(\mathbf{p}, \mathbf{h}) = \sum_{i=1}^{33} n_i (\mathbf{p} + 0.298)^{I_i} (\mathbf{h} - 0.720)^{J_i}, \quad (3)$$

where $\mathbf{q} = T/T^*$, $\mathbf{p} = p/p^*$, and $\mathbf{h} = h/h^*$ with $T^* = 860$ K, $p^* = 100$ MPa, and $h^* = 2800$ kJ kg⁻¹. The coefficients n_i and exponents I_i and J_i of Eq. (3) are listed in Table 5.

Computer-program verification

To assist the user in computer-program verification of Eqs. (2) and (3), Table 6 contains test values for calculated temperatures.

Table 4. Coefficients and exponents of the backward equation $T_{3a}(p, h)$ for subregion 3a in its dimensionless form, Eq. (2)

| i | I_i | J_i | n_i | i | I_i | J_i | n_i |
|-----|-------|-------|---|-----|-------|-------|---|
| 1 | -12 | 0 | -0.133 645 667 811 215 $\times 10^{-6}$ | 17 | -3 | 0 | -0.384 460 997 596 657 $\times 10^{-5}$ |
| 2 | -12 | 1 | 0.455 912 656 802 978 $\times 10^{-5}$ | 18 | -2 | 1 | 0.337 423 807 911 655 $\times 10^{-2}$ |
| 3 | -12 | 2 | -0.146 294 640 700 979 $\times 10^{-4}$ | 19 | -2 | 3 | -0.551 624 873 066 791 |
| 4 | -12 | 6 | 0.639 341 312 970 080 $\times 10^{-2}$ | 20 | -2 | 4 | 0.729 202 277 107 470 |
| 5 | -12 | 14 | 0.372 783 927 268 847 $\times 10^3$ | 21 | -1 | 0 | -0.992 522 757 376 041 $\times 10^{-2}$ |
| 6 | -12 | 16 | -0.718 654 377 460 447 $\times 10^4$ | 22 | -1 | 2 | -0.119 308 831 407 288 |
| 7 | -12 | 20 | 0.573 494 752 103 400 $\times 10^6$ | 23 | 0 | 0 | 0.793 929 190 615 421 |
| 8 | -12 | 22 | -0.267 569 329 111 439 $\times 10^7$ | 24 | 0 | 1 | 0.454 270 731 799 386 |
| 9 | -10 | 1 | -0.334 066 283 302 614 $\times 10^{-4}$ | 25 | 1 | 1 | 0.209 998 591 259 910 |
| 10 | -10 | 5 | -0.245 479 214 069 597 $\times 10^{-1}$ | 26 | 3 | 0 | -0.642 109 823 904 738 $\times 10^{-2}$ |
| 11 | -10 | 12 | 0.478 087 847 764 996 $\times 10^2$ | 27 | 3 | 1 | -0.235 155 868 604 540 $\times 10^{-1}$ |
| 12 | -8 | 0 | 0.764 664 131 818 904 $\times 10^{-5}$ | 28 | 4 | 0 | 0.252 233 108 341 612 $\times 10^{-2}$ |
| 13 | -8 | 2 | 0.128 350 627 676 972 $\times 10^{-2}$ | 29 | 4 | 3 | -0.764 885 133 368 119 $\times 10^{-2}$ |
| 14 | -8 | 4 | 0.171 219 081 377 331 $\times 10^{-1}$ | 30 | 10 | 4 | 0.136 176 427 574 291 $\times 10^{-1}$ |
| 15 | -8 | 10 | -0.851 007 304 583 213 $\times 10^1$ | 31 | 12 | 5 | -0.133 027 883 575 669 $\times 10^{-1}$ |
| 16 | -5 | 2 | -0.136 513 461 629 781 $\times 10^{-1}$ | | | | |

Table 5. Coefficients and exponents of the backward equation $T_{3b}(p, h)$ for subregion 3b in its dimensionless form, Eq. (3)

| i | I_i | J_i | n_i | i | I_i | J_i | n_i |
|-----|-------|-------|---|-----|-------|-------|---|
| 1 | -12 | 0 | 0.323 254 573 644 920 $\times 10^{-4}$ | 18 | -3 | 5 | -0.307 622 221 350 501 $\times 10^1$ |
| 2 | -12 | 1 | -0.127 575 556 587 181 $\times 10^{-3}$ | 19 | -2 | 0 | -0.574 011 959 864 879 $\times 10^{-1}$ |
| 3 | -10 | 0 | -0.475 851 877 356 068 $\times 10^{-3}$ | 20 | -2 | 4 | 0.503 471 360 939 849 $\times 10^1$ |
| 4 | -10 | 1 | 0.156 183 014 181 602 $\times 10^{-2}$ | 21 | -1 | 2 | -0.925 081 888 584 834 |
| 5 | -10 | 5 | 0.105 724 860 113 781 | 22 | -1 | 4 | 0.391 733 882 917 546 $\times 10^1$ |
| 6 | -10 | 10 | -0.858 514 221 132 534 $\times 10^2$ | 23 | -1 | 6 | -0.773 146 007 130 190 $\times 10^2$ |
| 7 | -10 | 12 | 0.724 140 095 480 911 $\times 10^3$ | 24 | -1 | 10 | 0.949 308 762 098 587 $\times 10^4$ |
| 8 | -8 | 0 | 0.296 475 810 273 257 $\times 10^{-2}$ | 25 | -1 | 14 | -0.141 043 719 679 409 $\times 10^7$ |
| 9 | -8 | 1 | -0.592 721 983 365 988 $\times 10^{-2}$ | 26 | -1 | 16 | 0.849 166 230 819 026 $\times 10^7$ |
| 10 | -8 | 2 | -0.126 305 422 818 666 $\times 10^{-1}$ | 27 | 0 | 0 | 0.861 095 729 446 704 |
| 11 | -8 | 4 | -0.115 716 196 364 853 | 28 | 0 | 2 | 0.323 346 442 811 720 |
| 12 | -8 | 10 | 0.849 000 969 739 595 $\times 10^2$ | 29 | 1 | 1 | 0.873 281 936 020 439 |
| 13 | -6 | 0 | -0.108 602 260 086 615 $\times 10^{-1}$ | 30 | 3 | 1 | -0.436 653 048 526 683 |
| 14 | -6 | 1 | 0.154 304 475 328 851 $\times 10^{-1}$ | 31 | 5 | 1 | 0.286 596 714 529 479 |
| 15 | -6 | 2 | 0.750 455 441 524 466 $\times 10^{-1}$ | 32 | 6 | 1 | -0.131 778 331 276 228 |
| 16 | -4 | 0 | 0.252 520 973 612 982 $\times 10^{-1}$ | 33 | 8 | 1 | 0.676 682 064 330 275 $\times 10^{-2}$ |
| 17 | -4 | 1 | -0.602 507 901 232 996 $\times 10^{-1}$ | | | | |

Table 6. Selected temperature values calculated from Eqs. (2) and (3) ^a

| Equation | p / MPa | h / kJ kg ⁻¹ | T / K |
|--------------------------|-----------|---------------------------|-------------------------------|
| $T_{3a}(p, h)$, Eq. (2) | 20 | 1700 | $6.293\ 083\ 892 \times 10^2$ |
| | 50 | 2000 | $6.905\ 718\ 338 \times 10^2$ |
| | 100 | 2100 | $7.336\ 163\ 014 \times 10^2$ |
| $T_{3b}(p, h)$, Eq. (3) | 20 | 2500 | $6.418\ 418\ 053 \times 10^2$ |
| | 50 | 2400 | $7.351\ 848\ 618 \times 10^2$ |
| | 100 | 2700 | $8.420\ 460\ 876 \times 10^2$ |

^a It is recommended that programmed functions be verified using 8 byte real values for all variables.

5.2 The Equations $v(p, h)$

The backward equation $v_{3a}(p, h)$ for subregion 3a has the following dimensionless form:

$$\frac{v_{3a}(p, h)}{v^*} = \mathbf{w}_{3a}(\mathbf{p}, \mathbf{h}) = \sum_{i=1}^{32} n_i (\mathbf{p} + 0.128)^{I_i} (\mathbf{h} - 0.727)^{J_i}, \quad (4)$$

where $\mathbf{w} = v/v^*$, $\mathbf{p} = p/p^*$, and $\mathbf{h} = h/h^*$ with $v^* = 0.0028 \text{ m}^3 \text{ kg}^{-1}$, $p^* = 100 \text{ MPa}$, and $h^* = 2100 \text{ kJ kg}^{-1}$. The coefficients n_i and exponents I_i and J_i of Eq. (4) are listed in Table 7.

The backward equation $v_{3b}(p, h)$ for subregion 3b reads in its dimensionless form

$$\frac{v_{3b}(p, h)}{v^*} = \mathbf{w}_{3b}(\mathbf{p}, \mathbf{h}) = \sum_{i=1}^{30} n_i (\mathbf{p} + 0.0661)^{I_i} (\mathbf{h} - 0.720)^{J_i}, \quad (5)$$

where $\mathbf{w} = v/v^*$, $\mathbf{p} = p/p^*$, and $\mathbf{h} = h/h^*$ with $v^* = 0.0088 \text{ m}^3 \text{ kg}^{-1}$, $p^* = 100 \text{ MPa}$, and $h^* = 2800 \text{ kJ kg}^{-1}$. The coefficients n_i and exponents I_i and J_i of Eq. (5) are listed in Table 8.

Computer-program verification

To assist the user in computer-program verification of Eqs. (4) and (5), Table 9 contains test values for calculated specific volumes.

Table 7. Coefficients and exponents of the backward equation $v_{3a}(p, h)$ for subregion 3a in its dimensionless form, Eq. (4)

| i | I_i | J_i | n_i | i | I_i | J_i | n_i |
|-----|-------|-------|---|-----|-------|-------|---|
| 1 | -12 | 6 | $0.529\ 944\ 062\ 966\ 028 \times 10^{-2}$ | 17 | -2 | 16 | $0.568\ 366\ 875\ 815\ 960 \times 10^4$ |
| 2 | -12 | 8 | $-0.170\ 099\ 690\ 234\ 461$ | 18 | -1 | 0 | $0.808\ 169\ 540\ 124\ 668 \times 10^{-2}$ |
| 3 | -12 | 12 | $0.111\ 323\ 814\ 312\ 927 \times 10^2$ | 19 | -1 | 1 | $0.172\ 416\ 341\ 519\ 307$ |
| 4 | -12 | 18 | $-0.217\ 898\ 123\ 145\ 125 \times 10^4$ | 20 | -1 | 2 | $0.104\ 270\ 175\ 292\ 927 \times 10^1$ |
| 5 | -10 | 4 | $-0.506\ 061\ 827\ 980\ 875 \times 10^{-3}$ | 21 | -1 | 3 | $-0.297\ 691\ 372\ 792\ 847$ |
| 6 | -10 | 7 | $0.556\ 495\ 239\ 685\ 324$ | 22 | 0 | 0 | $0.560\ 394\ 465\ 163\ 593$ |
| 7 | -10 | 10 | $-0.943\ 672\ 726\ 094\ 016 \times 10^1$ | 23 | 0 | 1 | $0.275\ 234\ 661\ 176\ 914$ |
| 8 | -8 | 5 | $-0.297\ 856\ 807\ 561\ 527$ | 24 | 1 | 0 | $-0.148\ 347\ 894\ 866\ 012$ |
| 9 | -8 | 12 | $0.939\ 353\ 943\ 717\ 186 \times 10^2$ | 25 | 1 | 1 | $-0.651\ 142\ 513\ 478\ 515 \times 10^{-1}$ |
| 10 | -6 | 3 | $0.192\ 944\ 939\ 465\ 981 \times 10^{-1}$ | 26 | 1 | 2 | $-0.292\ 468\ 715\ 386\ 302 \times 10^1$ |
| 11 | -6 | 4 | $0.421\ 740\ 664\ 704\ 763$ | 27 | 2 | 0 | $0.664\ 876\ 096\ 952\ 665 \times 10^{-1}$ |
| 12 | -6 | 22 | $-0.368\ 914\ 126\ 282\ 330 \times 10^7$ | 28 | 2 | 2 | $0.352\ 335\ 014\ 263\ 844 \times 10^1$ |
| 13 | -4 | 2 | $-0.737\ 566\ 847\ 600\ 639 \times 10^{-2}$ | 29 | 3 | 0 | $-0.146\ 340\ 792\ 313\ 332 \times 10^{-1}$ |
| 14 | -4 | 3 | $-0.354\ 753\ 242\ 424\ 366$ | 30 | 4 | 2 | $-0.224\ 503\ 486\ 668\ 184 \times 10^1$ |
| 15 | -3 | 7 | $-0.199\ 768\ 169\ 338\ 727 \times 10^1$ | 31 | 5 | 2 | $0.110\ 533\ 464\ 706\ 142 \times 10^1$ |
| 16 | -2 | 3 | $0.115\ 456\ 297\ 059\ 049 \times 10^1$ | 32 | 8 | 2 | $-0.408\ 757\ 344\ 495\ 612 \times 10^{-1}$ |

Table 8. Coefficients and exponents of the backward equation $v_{3b}(p, h)$ for subregion 3b in its dimensionless form, Eq. (5)

| i | I_i | J_i | n_i | i | I_i | J_i | n_i |
|-----|-------|-------|---|-----|-------|-------|---|
| 1 | -12 | 0 | $-0.225\ 196\ 934\ 336\ 318 \times 10^{-8}$ | 16 | -4 | 6 | $-0.321\ 087\ 965\ 668\ 917 \times 10^1$ |
| 2 | -12 | 1 | $0.140\ 674\ 363\ 313\ 486 \times 10^{-7}$ | 17 | -4 | 10 | $0.607\ 567\ 815\ 637\ 771 \times 10^3$ |
| 3 | -8 | 0 | $0.233\ 784\ 085\ 280\ 560 \times 10^{-5}$ | 18 | -3 | 0 | $0.557\ 686\ 450\ 685\ 932 \times 10^{-3}$ |
| 4 | -8 | 1 | $-0.331\ 833\ 715\ 229\ 001 \times 10^{-4}$ | 19 | -3 | 2 | $0.187\ 499\ 040\ 029\ 550$ |
| 5 | -8 | 3 | $0.107\ 956\ 778\ 514\ 318 \times 10^{-2}$ | 20 | -2 | 1 | $0.905\ 368\ 030\ 448\ 107 \times 10^{-2}$ |
| 6 | -8 | 6 | $-0.271\ 382\ 067\ 378\ 863$ | 21 | -2 | 2 | $0.285\ 417\ 173\ 048\ 685$ |
| 7 | -8 | 7 | $0.107\ 202\ 262\ 490\ 333 \times 10^1$ | 22 | -1 | 0 | $0.329\ 924\ 030\ 996\ 098 \times 10^{-1}$ |
| 8 | -8 | 8 | $-0.853\ 821\ 329\ 075\ 382$ | 23 | -1 | 1 | $0.239\ 897\ 419\ 685\ 483$ |
| 9 | -6 | 0 | $-0.215\ 214\ 194\ 340\ 526 \times 10^{-4}$ | 24 | -1 | 4 | $0.482\ 754\ 995\ 951\ 394 \times 10^1$ |
| 10 | -6 | 1 | $0.769\ 656\ 088\ 222\ 730 \times 10^{-3}$ | 25 | -1 | 5 | $-0.118\ 035\ 753\ 702\ 231 \times 10^2$ |
| 11 | -6 | 2 | $-0.431\ 136\ 580\ 433\ 864 \times 10^{-2}$ | 26 | 0 | 0 | $0.169\ 490\ 044\ 091\ 791$ |
| 12 | -6 | 5 | $0.453\ 342\ 167\ 309\ 331$ | 27 | 1 | 0 | $-0.179\ 967\ 222\ 507\ 787 \times 10^{-1}$ |
| 13 | -6 | 6 | $-0.507\ 749\ 535\ 873\ 652$ | 28 | 1 | 1 | $0.371\ 810\ 116\ 332\ 674 \times 10^{-1}$ |
| 14 | -6 | 10 | $-0.100\ 475\ 154\ 528\ 389 \times 10^3$ | 29 | 2 | 2 | $-0.536\ 288\ 335\ 065\ 096 \times 10^{-1}$ |
| 15 | -4 | 3 | $-0.219\ 201\ 924\ 648\ 793$ | 30 | 2 | 6 | $0.160\ 697\ 101\ 092\ 520 \times 10^1$ |

Table 9. Selected specific volume values calculated from Eqs. (4) and (5) ^a

| Equation | p / MPa | $h / \text{kJ kg}^{-1}$ | $v / \text{m}^3 \text{kg}^{-1}$ |
|--------------------------|------------------|-------------------------|----------------------------------|
| $v_{3a}(p, h)$, Eq. (4) | 20 | 1700 | $1.749\,903\,962 \times 10^{-3}$ |
| | 50 | 2000 | $1.908\,139\,035 \times 10^{-3}$ |
| | 100 | 2100 | $1.676\,229\,776 \times 10^{-3}$ |
| $v_{3b}(p, h)$, Eq. (5) | 20 | 2500 | $6.670\,547\,043 \times 10^{-3}$ |
| | 50 | 2400 | $2.801\,244\,590 \times 10^{-3}$ |
| | 100 | 2700 | $2.404\,234\,998 \times 10^{-3}$ |

^a It is recommended that programmed functions be verified using 8 byte real values for all variables.

5.3 Numerical Consistency with the Basic Equation of IAPWS-IF97

The maximum temperature differences and related root-mean-square differences between the calculated temperature Eqs. (2) and (3) and the IAPWS-IF97 basic equation $f_3^{97}(v, T)$ in comparison with the permissible differences are listed in Table 10. The calculation of the root-mean-square values is described in Section 1.

Table 10 also contains the maximum relative deviations and root-mean-square relative deviations for specific volume of Eqs. (4) and (5) from IAPWS-IF97.

The critical temperature and the critical volume are met exactly by the equations $T(p, h)$ and $v(p, h)$.

Table 10. Maximum differences and root-mean-square differences of the temperature calculated from Eqs. (2) and (3) and specific volume calculated from Eqs. (4) and (5) to the IAPWS-IF97 basic equation $f_3^{97}(v, T)$ and related permissible values

| Subregion | Equation | $ \Delta T _{\text{tol}}$ | $ \Delta T _{\text{max}}$ | $ \Delta T _{\text{RMS}}$ |
|-----------|----------|-----------------------------|-----------------------------|-----------------------------|
| 3a | (2) | 25 mK | 23.6 mK | 10.5 mK |
| 3b | (3) | 25 mK | 19.6 mK | 9.6 mK |
| Subregion | Equation | $ \Delta v/v _{\text{tol}}$ | $ \Delta v/v _{\text{max}}$ | $ \Delta v/v _{\text{RMS}}$ |
| 3a | (4) | 0.01 % | 0.0080 % | 0.0032 % |
| 3b | (5) | 0.01 % | 0.0095 % | 0.0042 % |

5.4 Consistency at Boundary Between Subregions

The maximum temperature difference between the two backward equations, Eq. (2) and Eq. (3), along the boundary $h_{3ab}(p)$, Eq. (1), has the following value

$$|\Delta T|_{\text{max}} = \left| T_{3a}(p, h_{3ab}(p)) - T_{3b}(p, h_{3ab}(p)) \right|_{\text{max}} = 0.37 \text{ mK}.$$

Thus, the temperature differences between the two backward functions $T(p, h)$ of the adjacent subregions are smaller than the numerical consistencies with the IAPWS-IF97 equations.

The relative specific volume differences between the two backward equations $v(p, h)$ of the adjacent subregions 3a and 3b are also smaller than the numerical consistencies of these equations with the IAPWS-IF97 basic equation. Along the boundary $h_{3ab}(p)$, Eq. (1), the maximum difference between the corresponding equations was determined as:

$$\left| \frac{\Delta v}{v} \right|_{\max} = \left| \frac{v_{3a}(p, h_{3ab}(p)) - v_{3b}(p, h_{3ab}(p))}{v_{3b}(p, h_{3ab}(p))} \right|_{\max} = 0.00015\% .$$

6 Backward Equations $T(p, s)$ and $v(p, s)$ for Subregions 3a and 3b

6.1 The Equations $T(p, s)$

The backward equation $T_{3a}(p, s)$ for subregion 3a has the following dimensionless form:

$$\frac{T_{3a}(p, s)}{T^*} = \mathbf{q}_{3a}(\mathbf{p}, \mathbf{s}) = \sum_{i=1}^{33} n_i (\mathbf{p} + 0.240)^{I_i} (\mathbf{s} - 0.703)^{J_i} , \quad (6)$$

where $\mathbf{q} = T/T^*$, $\mathbf{p} = p/p^*$, and $\mathbf{s} = s/s^*$ with $T^* = 760$ K, $p^* = 100$ MPa, and $s^* = 4.4$ kJ kg⁻¹ K⁻¹. The coefficients n_i and exponents I_i and J_i of Eq. (6) are listed in Table 11.

The backward equation $T_{3b}(p, s)$ for subregion 3b reads in its dimensionless form

$$\frac{T_{3b}(p, s)}{T^*} = \mathbf{q}_{3b}(\mathbf{p}, \mathbf{s}) = \sum_{i=1}^{28} n_i (\mathbf{p} + 0.760)^{I_i} (\mathbf{s} - 0.818)^{J_i} , \quad (7)$$

where $\mathbf{q} = T/T^*$, $\mathbf{p} = p/p^*$, and $\mathbf{s} = s/s^*$ with $T^* = 860$ K, $p^* = 100$ MPa, and $s^* = 5.3$ kJ kg⁻¹ K⁻¹. The coefficients n_i and exponents I_i and J_i of Eq. (7) are listed in Table 12.

Computer-program verification

To assist the user in computer-program verification of Eqs. (6) and (7), Table 13 contains test values for calculated temperatures.

Table 11. Coefficients and exponents of the backward equation $T_{3a}(p, s)$ for subregion 3a in its dimensionless form, Eq. (6)

| i | I_i | J_i | n_i | i | I_i | J_i | n_i |
|-----|-------|-------|---|-----|-------|-------|---|
| 1 | -12 | 28 | $0.150\ 042\ 008\ 263\ 875 \times 10^{10}$ | 18 | -4 | 10 | $-0.368\ 275\ 545\ 889\ 071 \times 10^3$ |
| 2 | -12 | 32 | $-0.159\ 397\ 258\ 480\ 424 \times 10^{12}$ | 19 | -4 | 36 | $0.664\ 768\ 904\ 779\ 177 \times 10^{16}$ |
| 3 | -10 | 4 | $0.502\ 181\ 140\ 217\ 975 \times 10^{-3}$ | 20 | -2 | 1 | $0.449\ 359\ 251\ 958\ 880 \times 10^{-1}$ |
| 4 | -10 | 10 | $-0.672\ 057\ 767\ 855\ 466 \times 10^2$ | 21 | -2 | 4 | $-0.422\ 897\ 836\ 099\ 655 \times 10^1$ |
| 5 | -10 | 12 | $0.145\ 058\ 545\ 404\ 456 \times 10^4$ | 22 | -1 | 1 | $-0.240\ 614\ 376\ 434\ 179$ |
| 6 | -10 | 14 | $-0.823\ 889\ 534\ 888\ 890 \times 10^4$ | 23 | -1 | 6 | $-0.474\ 341\ 365\ 254\ 924 \times 10^1$ |
| 7 | -8 | 5 | $-0.154\ 852\ 214\ 233\ 853$ | 24 | 0 | 0 | $0.724\ 093\ 999\ 126\ 110$ |
| 8 | -8 | 7 | $0.112\ 305\ 046\ 746\ 695 \times 10^2$ | 25 | 0 | 1 | $0.923\ 874\ 349\ 695\ 897$ |
| 9 | -8 | 8 | $-0.297\ 000\ 213\ 482\ 822 \times 10^2$ | 26 | 0 | 4 | $0.399\ 043\ 655\ 281\ 015 \times 10^1$ |
| 10 | -8 | 28 | $0.438\ 565\ 132\ 635\ 495 \times 10^{11}$ | 27 | 1 | 0 | $0.384\ 066\ 651\ 868\ 009 \times 10^{-1}$ |
| 11 | -6 | 2 | $0.137\ 837\ 838\ 635\ 464 \times 10^{-2}$ | 28 | 2 | 0 | $-0.359\ 344\ 365\ 571\ 848 \times 10^{-2}$ |
| 12 | -6 | 6 | $-0.297\ 478\ 527\ 157\ 462 \times 10^1$ | 29 | 2 | 3 | $-0.735\ 196\ 448\ 821\ 653$ |
| 13 | -6 | 32 | $0.971\ 777\ 947\ 349\ 413 \times 10^{13}$ | 30 | 3 | 2 | $0.188\ 367\ 048\ 396\ 131$ |
| 14 | -5 | 0 | $-0.571\ 527\ 767\ 052\ 398 \times 10^{-4}$ | 31 | 8 | 0 | $0.141\ 064\ 266\ 818\ 704 \times 10^{-3}$ |
| 15 | -5 | 14 | $0.288\ 307\ 949\ 778\ 420 \times 10^5$ | 32 | 8 | 1 | $-0.257\ 418\ 501\ 496\ 337 \times 10^{-2}$ |
| 16 | -5 | 32 | $-0.744\ 428\ 289\ 262\ 703 \times 10^{14}$ | 33 | 10 | 2 | $0.123\ 220\ 024\ 851\ 555 \times 10^{-2}$ |
| 17 | -4 | 6 | $0.128\ 017\ 324\ 848\ 921 \times 10^2$ | | | | |

Table 12. Coefficients and exponents of the backward equation $T_{3b}(p, s)$ for subregion 3b in its dimensionless form, Eq. (7)

| i | I_i | J_i | n_i | i | I_i | J_i | n_i |
|-----|-------|-------|--|-----|-------|-------|---|
| 1 | -12 | 1 | $0.527\ 111\ 701\ 601\ 660$ | 15 | -5 | 6 | $0.880\ 531\ 517\ 490\ 555 \times 10^3$ |
| 2 | -12 | 3 | $-0.401\ 317\ 830\ 052\ 742 \times 10^2$ | 16 | -4 | 12 | $0.265\ 015\ 592\ 794\ 626 \times 10^7$ |
| 3 | -12 | 4 | $0.153\ 020\ 073\ 134\ 484 \times 10^3$ | 17 | -3 | 1 | $-0.359\ 287\ 150\ 025\ 783$ |
| 4 | -12 | 7 | $-0.224\ 799\ 398\ 218\ 827 \times 10^4$ | 18 | -3 | 6 | $-0.656\ 991\ 567\ 673\ 753 \times 10^3$ |
| 5 | -8 | 0 | $-0.193\ 993\ 484\ 669\ 048$ | 19 | -2 | 2 | $0.241\ 768\ 149\ 185\ 367 \times 10^1$ |
| 6 | -8 | 1 | $-0.140\ 467\ 557\ 893\ 768 \times 10^1$ | 20 | 0 | 0 | $0.856\ 873\ 461\ 222\ 588$ |
| 7 | -8 | 3 | $0.426\ 799\ 878\ 114\ 024 \times 10^2$ | 21 | 2 | 1 | $0.655\ 143\ 675\ 313\ 458$ |
| 8 | -6 | 0 | $0.752\ 810\ 643\ 416\ 743$ | 22 | 3 | 1 | $-0.213\ 535\ 213\ 206\ 406$ |
| 9 | -6 | 2 | $0.226\ 657\ 238\ 616\ 417 \times 10^2$ | 23 | 4 | 0 | $0.562\ 974\ 957\ 606\ 348 \times 10^{-2}$ |
| 10 | -6 | 4 | $-0.622\ 873\ 556\ 909\ 932 \times 10^3$ | 24 | 5 | 24 | $-0.316\ 955\ 725\ 450\ 471 \times 10^{15}$ |
| 11 | -5 | 0 | $-0.660\ 823\ 667\ 935\ 396$ | 25 | 6 | 0 | $-0.699\ 997\ 000\ 152\ 457 \times 10^{-3}$ |
| 12 | -5 | 1 | $0.841\ 267\ 087\ 271\ 658$ | 26 | 8 | 3 | $0.119\ 845\ 803\ 210\ 767 \times 10^{-1}$ |
| 13 | -5 | 2 | $-0.253\ 717\ 501\ 764\ 397 \times 10^2$ | 27 | 12 | 1 | $0.193\ 848\ 122\ 022\ 095 \times 10^{-4}$ |
| 14 | -5 | 4 | $0.485\ 708\ 963\ 532\ 948 \times 10^3$ | 28 | 14 | 2 | $-0.215\ 095\ 749\ 182\ 309 \times 10^{-4}$ |

Table 13. Selected temperature values calculated from Eqs. (6) and (7)^a

| Equation | p / MPa | $s / \text{kJ kg}^{-1} \text{K}^{-1}$ | T / K |
|--------------------------|------------------|---------------------------------------|-------------------------------|
| $T_{3a}(p, s)$, Eq. (6) | 20 | 3.7 | $6.208\ 841\ 563 \times 10^2$ |
| | 50 | 3.5 | $6.181\ 549\ 029 \times 10^2$ |
| | 100 | 4.0 | $7.056\ 880\ 237 \times 10^2$ |
| $T_{3b}(p, s)$, Eq. (7) | 20 | 5.0 | $6.401\ 176\ 443 \times 10^2$ |
| | 50 | 4.5 | $7.163\ 687\ 517 \times 10^2$ |
| | 100 | 5.0 | $8.474\ 332\ 825 \times 10^2$ |

^a It is recommended that programmed functions be verified using 8 byte real values for all variables.

6.2 The Equations $v(p, s)$

The backward equation $v_{3a}(p, s)$ for subregion 3a has the following dimensionless form:

$$\frac{v_{3a}(p, s)}{v^*} = \mathbf{w}_{3a}(\mathbf{p}, \mathbf{s}) = \sum_{i=1}^{28} n_i (\mathbf{p} + 0.187)^{I_i} (\mathbf{s} - 0.755)^{J_i}, \quad (8)$$

where $\mathbf{w} = v/v^*$, $\mathbf{p} = p/p^*$, and $\mathbf{s} = s/s^*$ with $v^* = 0.0028 \text{ m}^3 \text{ kg}^{-1}$, $p^* = 100 \text{ MPa}$, and $s^* = 4.4 \text{ kJ kg}^{-1} \text{ K}^{-1}$. The coefficients n_i and exponents I_i and J_i of Eq. (8) are listed in Table 14.

The backward equation $v_{3b}(p, s)$ for subregion 3b reads in its dimensionless form

$$\frac{v_{3b}(p, s)}{v^*} = \mathbf{w}_{3b}(\mathbf{p}, \mathbf{s}) = \sum_{i=1}^{31} n_i (\mathbf{p} + 0.298)^{I_i} (\mathbf{s} - 0.816)^{J_i}, \quad (9)$$

where $\mathbf{w} = v/v^*$, $\mathbf{p} = p/p^*$, and $\mathbf{s} = s/s^*$ with $v^* = 0.0088 \text{ m}^3 \text{ kg}^{-1}$, $p^* = 100 \text{ MPa}$, and $s^* = 5.3 \text{ kJ kg}^{-1} \text{ K}^{-1}$. The coefficients n_i and exponents I_i and J_i of Eq. (9) are listed in Table 15.

Computer-program verification

To assist the user in computer-program verification of Eqs. (8) and (9), Table 16 contains test values for calculated specific volumes.

Table 14. Coefficients and exponents of the backward equation $v_{3a}(p, s)$ for subregion 3a in its dimensionless form, Eq. (8).

| i | I_i | J_i | n_i | i | I_i | J_i | n_i |
|-----|-------|-------|---|-----|-------|-------|---|
| 1 | -12 | 10 | $0.795\ 544\ 074\ 093\ 975 \times 10^2$ | 15 | -3 | 2 | $-0.118\ 008\ 384\ 666\ 987$ |
| 2 | -12 | 12 | $-0.238\ 261\ 242\ 984\ 590 \times 10^4$ | 16 | -3 | 4 | $0.253\ 798\ 642\ 355\ 900 \times 10^1$ |
| 3 | -12 | 14 | $0.176\ 813\ 100\ 617\ 787 \times 10^5$ | 17 | -2 | 3 | $0.965\ 127\ 704\ 669\ 424$ |
| 4 | -10 | 4 | $-0.110\ 524\ 727\ 080\ 379 \times 10^{-2}$ | 18 | -2 | 8 | $-0.282\ 172\ 420\ 532\ 826 \times 10^2$ |
| 5 | -10 | 8 | $-0.153\ 213\ 833\ 655\ 326 \times 10^2$ | 19 | -1 | 1 | $0.203\ 224\ 612\ 353\ 823$ |
| 6 | -10 | 10 | $0.297\ 544\ 599\ 376\ 982 \times 10^3$ | 20 | -1 | 2 | $0.110\ 648\ 186\ 063\ 513 \times 10^1$ |
| 7 | -10 | 20 | $-0.350\ 315\ 206\ 871\ 242 \times 10^8$ | 21 | 0 | 0 | $0.526\ 127\ 948\ 451\ 280$ |
| 8 | -8 | 5 | $0.277\ 513\ 761\ 062\ 119$ | 22 | 0 | 1 | $0.277\ 000\ 018\ 736\ 321$ |
| 9 | -8 | 6 | $-0.523\ 964\ 271\ 036\ 888$ | 23 | 0 | 3 | $0.108\ 153\ 340\ 501\ 132 \times 10^1$ |
| 10 | -8 | 14 | $-0.148\ 011\ 182\ 995\ 403 \times 10^6$ | 24 | 1 | 0 | $-0.744\ 127\ 885\ 357\ 893 \times 10^{-1}$ |
| 11 | -8 | 16 | $0.160\ 014\ 899\ 374\ 266 \times 10^7$ | 25 | 2 | 0 | $0.164\ 094\ 443\ 541\ 384 \times 10^{-1}$ |
| 12 | -6 | 28 | $0.170\ 802\ 322\ 663\ 427 \times 10^{13}$ | 26 | 4 | 2 | $-0.680\ 468\ 275\ 301\ 065 \times 10^{-1}$ |
| 13 | -5 | 1 | $0.246\ 866\ 996\ 006\ 494 \times 10^{-3}$ | 27 | 5 | 2 | $0.257\ 988\ 576\ 101\ 640 \times 10^{-1}$ |
| 14 | -4 | 5 | $0.165\ 326\ 084\ 797\ 980 \times 10^1$ | 28 | 6 | 0 | $-0.145\ 749\ 861\ 944\ 416 \times 10^{-3}$ |

Table 15. Coefficients and exponents of the backward equation $v_{3b}(p, s)$ for subregion 3b in its dimensionless form, Eq. (9)

| i | I_i | J_i | n_i | i | I_i | J_i | n_i |
|-----|-------|-------|---|-----|-------|-------|---|
| 1 | -12 | 0 | $0.591\ 599\ 780\ 322\ 238 \times 10^{-4}$ | 17 | -4 | 2 | $-0.121\ 613\ 320\ 606\ 788 \times 10^2$ |
| 2 | -12 | 1 | $-0.185\ 465\ 997\ 137\ 856 \times 10^{-2}$ | 18 | -4 | 3 | $0.167\ 637\ 540\ 957\ 944 \times 10^1$ |
| 3 | -12 | 2 | $0.104\ 190\ 510\ 480\ 013 \times 10^{-1}$ | 19 | -3 | 1 | $-0.744\ 135\ 838\ 773\ 463 \times 10^1$ |
| 4 | -12 | 3 | $0.598\ 647\ 302\ 038\ 590 \times 10^{-2}$ | 20 | -2 | 0 | $0.378\ 168\ 091\ 437\ 659 \times 10^{-1}$ |
| 5 | -12 | 5 | $-0.771\ 391\ 189\ 901\ 699$ | 21 | -2 | 1 | $0.401\ 432\ 203\ 027\ 688 \times 10^1$ |
| 6 | -12 | 6 | $0.172\ 549\ 765\ 557\ 036 \times 10^1$ | 22 | -2 | 2 | $0.160\ 279\ 837\ 479\ 185 \times 10^2$ |
| 7 | -10 | 0 | $-0.467\ 076\ 079\ 846\ 526 \times 10^{-3}$ | 23 | -2 | 3 | $0.317\ 848\ 779\ 347\ 728 \times 10^1$ |
| 8 | -10 | 1 | $0.134\ 533\ 823\ 384\ 439 \times 10^{-1}$ | 24 | -2 | 4 | $-0.358\ 362\ 310\ 304\ 853 \times 10^1$ |
| 9 | -10 | 2 | $-0.808\ 094\ 336\ 805\ 495 \times 10^{-1}$ | 25 | -2 | 12 | $-0.115\ 995\ 260\ 446\ 827 \times 10^7$ |
| 10 | -10 | 4 | $0.508\ 139\ 374\ 365\ 767$ | 26 | 0 | 0 | $0.199\ 256\ 573\ 577\ 909$ |
| 11 | -8 | 0 | $0.128\ 584\ 643\ 361\ 683 \times 10^{-2}$ | 27 | 0 | 1 | $-0.122\ 270\ 624\ 794\ 624$ |
| 12 | -5 | 1 | $-0.163\ 899\ 353\ 915\ 435 \times 10^1$ | 28 | 0 | 2 | $-0.191\ 449\ 143\ 716\ 586 \times 10^2$ |
| 13 | -5 | 2 | $0.586\ 938\ 199\ 318\ 063 \times 10^1$ | 29 | 1 | 0 | $-0.150\ 448\ 002\ 905\ 284 \times 10^{-1}$ |
| 14 | -5 | 3 | $-0.292\ 466\ 667\ 918\ 613 \times 10^1$ | 30 | 1 | 2 | $0.146\ 407\ 900\ 162\ 154 \times 10^2$ |
| 15 | -4 | 0 | $-0.614\ 076\ 301\ 499\ 537 \times 10^{-2}$ | 31 | 2 | 2 | $-0.327\ 477\ 787\ 188\ 230 \times 10^1$ |
| 16 | -4 | 1 | $0.576\ 199\ 014\ 049\ 172 \times 10^1$ | | | | |

Table 16. Selected specific volume values calculated from Eqs. (8) and (9) ^a

| Equation | p / MPa | $s / \text{kJ kg}^{-1} \text{K}^{-1}$ | $v / \text{m}^3 \text{kg}^{-1}$ |
|--------------------------|------------------|---------------------------------------|----------------------------------|
| $v_{3a}(p, s)$, Eq. (8) | 20 | 3.7 | $1.639\,890\,984 \times 10^{-3}$ |
| | 50 | 3.5 | $1.423\,030\,205 \times 10^{-3}$ |
| | 100 | 4.0 | $1.555\,893\,131 \times 10^{-3}$ |
| $v_{3b}(p, s)$, Eq. (9) | 20 | 5.0 | $6.262\,101\,987 \times 10^{-3}$ |
| | 50 | 4.5 | $2.332\,634\,294 \times 10^{-3}$ |
| | 100 | 5.0 | $2.449\,610\,757 \times 10^{-3}$ |

^a It is recommended that programmed functions be verified using 8 byte real values for all variables.

6.3 Numerical Consistency with the Basic Equation of IAPWS-IF97

The maximum temperature differences and related root-mean-square differences between the temperatures calculated from Eqs. (6) and (7) and the IAPWS-IF97 basic equation $f_3^{97}(v, T)$ in comparison with the permissible differences are listed in Table 17.

Table 17 also contains the maximum relative deviations and root-mean-square relative deviations for the specific volume of Eqs. (8) and (9) from IAPWS-IF97.

The critical temperature and the critical volume are met exactly by the equations $T(p, s)$ and $v(p, s)$.

Table 17. Maximum differences and root-mean-square differences of the temperature calculated from Eqs. (6) and (7), and specific volume calculated from Eqs. (8) and (9) to the IAPWS-IF97 basic equation $f_3^{97}(v, T)$ and related permissible values

| Subregion | Equation | $ \Delta T _{\text{tol}}$ | $ \Delta T _{\text{max}}$ | $ \Delta T _{\text{RMS}}$ |
|-----------|----------|-----------------------------|-----------------------------|-----------------------------|
| 3a | (6) | 25 mK | 24.8 mK | 11.2 mK |
| 3b | (7) | 25 mK | 22.1 mK | 10.1 mK |
| Subregion | Equation | $ \Delta v/v _{\text{tol}}$ | $ \Delta v/v _{\text{max}}$ | $ \Delta v/v _{\text{RMS}}$ |
| 3a | (8) | 0.01 % | 0.0096 % | 0.0052 % |
| 3b | (9) | 0.01 % | 0.0077 % | 0.0037 % |

6.4 Consistency at Boundary Between Subregions

The maximum temperature difference between the two backward equations, Eq. (6) and Eq. (7), along the boundary s_c , has the following value

$$|\Delta T|_{\text{max}} = |T_{3a}(p, s_c) - T_{3b}(p, s_c)|_{\text{max}} = 0.093 \text{ mK}.$$

Thus, the temperature differences between the two backward functions $T(p,s)$ of the adjacent subregions are smaller than the numerical consistencies with the IAPWS-IF97 equations.

The relative specific volume differences between the two backward equations $v(p,s)$, Eqs. (8) and (9), of the adjacent subregions are also smaller than the numerical consistencies of these equations with the IAPWS-IF97 basic equation. Along the boundary s_c , the maximum difference between the corresponding equations was determined as

$$\left| \frac{\Delta v}{v} \right|_{\max} = \left| \frac{v_{3a}(p, s_c) - v_{3b}(p, s_c)}{v_{3b}(p, s_c)} \right|_{\max} = 0.00046\% .$$

7 Computing Time in Relation to IAPWS-IF97

A very important motivation for the development of the backward equations $T(p,h)$, $v(p,h)$ and $T(p,s)$, $v(p,s)$ for region 3 was reducing the computing time to obtain thermodynamic properties and differential quotients from given variables (p,h) and (p,s) . In IAPWS-IF97, time-consuming iterations, e.g., the 2-dimensional Newton method, are required. Using the $T_3(p,h)$, $v_3(p,h)$, $T_3(p,s)$ and $v_3(p,s)$ equations, the calculation speed is about 20 times faster than that of the 2-dimensional Newton method.

The numerical consistency of T and v obtained in this way is sufficient for most heat cycle calculations.

For users not satisfied with the numerical consistency of the backward equations, the equations are still recommended for generating starting points for the iterative process. They will significantly reduce the time required to reach the convergence criteria of the iteration.

8 References

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